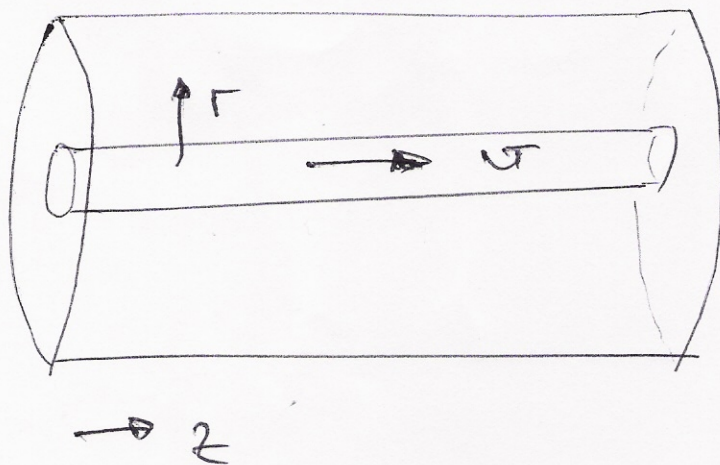


EXAMPLE SHEET VI

1)



Steady, ϕ & z - indep. B.C.

2 expect (try!) solution:

$$\underline{u} = (0, 0, \omega) = \omega(r) \underline{e}_z$$

$$\rho = \rho(z) \text{ (from } r \text{ \& } \phi \text{ momentum eqns)}$$

3 only z -momentum eqn is non zero & simplifies to:

$$\frac{dp}{dz} = \mu \nabla^2 \omega$$

$$\frac{1}{\mu} G = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \omega}{\partial r} \right) = \frac{\partial^2 \omega}{\partial r^2} + \frac{1}{r} \frac{\partial \omega}{\partial r}$$

$$\frac{q}{\mu} r = \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right)$$

$$\frac{1}{2} \frac{q}{\mu} r^2 + \tilde{A} = r \frac{\partial w}{\partial r}$$

$$\frac{1}{2} \frac{q}{\mu} r + \frac{\tilde{A}}{r} = \frac{\partial w}{\partial r}$$

$$w = \frac{1}{4} \frac{q}{\mu} r^2 + \tilde{A} \ln r + \tilde{B}$$

$$w = \frac{1}{4} \frac{q}{\mu} r^2 + A \ln \left(\frac{r}{b} \right) + B$$

B.C.:

$$w(r=b) = 0$$

$$B = -\frac{1}{4} \frac{q}{\mu} b^2$$

$$w(r=a) = \psi$$

$$\psi = \frac{1}{4} \frac{q}{\mu} (a^2 - b^2) + A \ln \left(\frac{a}{b} \right)$$

$$A = \frac{\psi + \frac{1}{4} \frac{q}{\mu} b^2}{\ln \left(\frac{a}{b} \right)} + \frac{\frac{1}{4} \frac{q}{\mu} (b^2 - a^2)}{\ln \left(\frac{a}{b} \right)}$$

$$\omega(r) = \frac{1}{4} \frac{q}{\mu} \left[(r^2 - b^2) + \frac{(b^2 - a^2) \ln\left(\frac{r}{b}\right)}{\ln\left(\frac{a}{b}\right)} \right] + \sigma \frac{\ln\left(\frac{r}{b}\right)}{\ln\left(\frac{a}{b}\right)}$$

Dref:

$$D = \underbrace{\tau_w}_{\text{wall shear}} \cdot \underbrace{2\pi a}_{\text{circumfer.}} \cdot \underbrace{L}_{\text{length.}}$$

walled end.

$$\tau_w = t_z$$

$$t_i = \tau_{ij} n_j$$

$$\underline{n} = -\underline{e}_r$$

$$\tau_{ij} = -p \delta_{ij} + 2\mu \epsilon_{ij}$$

$$i = 1, 2, 3 \rightsquigarrow i = r, \varphi, z$$

$$n_r = -1 \quad n_\varphi = 0 \quad n_z = 0$$

$$t_z = \tau_{zj} n_j$$

only $j=r$ survives

$$t_z = -2\mu \epsilon_{zr} = -\mu \left(\cancel{\frac{\partial v}{\partial z}} + \frac{\partial w}{\partial r} \right) = \mu \frac{\partial w}{\partial r}$$

$$\frac{\partial w}{\partial r} = \frac{1}{4} \frac{g}{\mu} \left[2r + \frac{b^2 - a^2}{r \ln(a/b)} \right] + \frac{v}{r \ln(a/b)}$$

$$\frac{D}{L} = 2\pi a \tau_w |_{r=a}$$

W
drop
per
unit
length

$$= - \left(\frac{\pi}{2} \frac{g a}{\mu} \left[2a + \frac{b^2 - a^2}{a \ln(a/b)} \right] + \frac{2\pi \mu v}{\ln(a/b)} \right)$$

(ii) $\frac{D}{L} = 0$ for

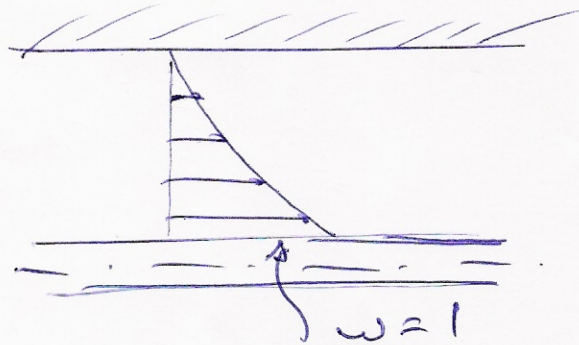
$$G = G_0 = - \frac{2\pi \mu v}{\ln(a/b) \frac{\pi}{2} a \left[2a + \frac{b^2 - a^2}{a \ln(a/b)} \right]}$$

$$G_0 = - \frac{4 \mu v}{a} \frac{1}{2a \ln(a/b) + \frac{b^2 - a^2}{a}}$$

$$G_0 = - \frac{4 \mu v}{a^2} \frac{1}{2 \ln(a/b) + \left(\frac{b}{a}\right)^2 - 1}$$

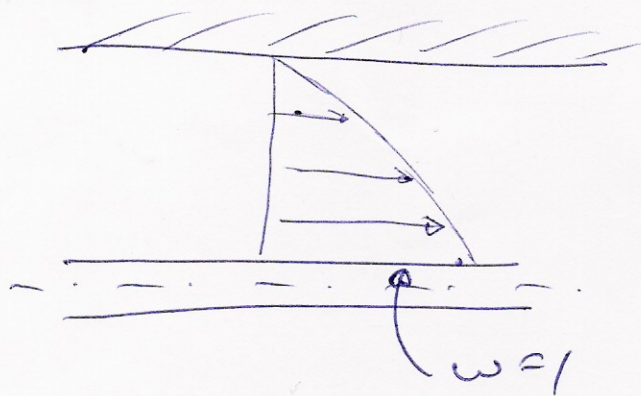
velocity profiles:

$Q = 0$



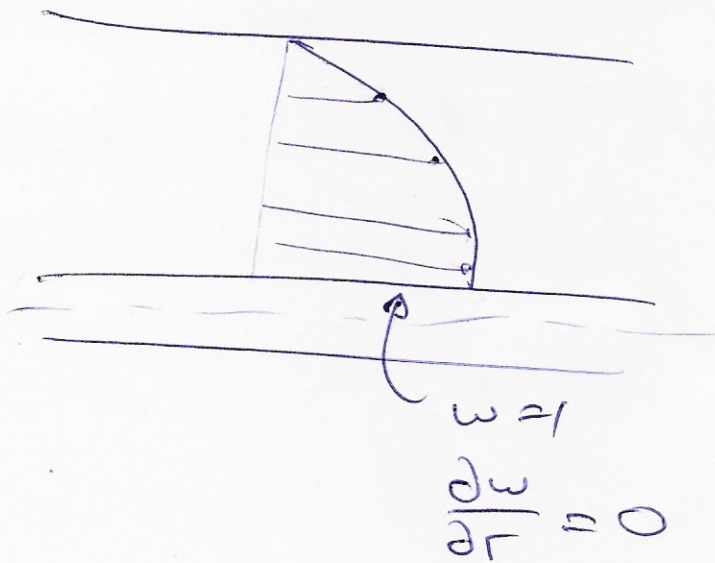
Flow only driven by boundary motion

$Q < 0$



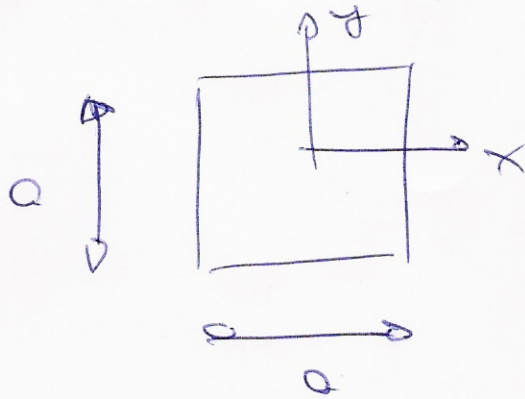
Combination of boundary & pressure driven flow [flow opposite press. gradient]

$Q = Q_0 < 0$



Critical case. No drag.

2) parallel steady flow



$$\frac{G}{\mu} = \nabla^2 \omega$$

$$\omega|_{x=\pm a/2} = \omega|_{y=\pm a/2} = 0$$

$$\omega = \omega_p(y) + \omega_H(x, y)$$

Choose ω_p s.t. it fulfills the BC.

$$\omega_p(y) = \frac{1}{2} \frac{G}{\mu} \left(y^2 - \left(\frac{a}{2} \right)^2 \right)$$

$$\nabla^2 \omega_p = \frac{G}{\mu} \quad \& \quad \omega_p(y = \pm a/2) = 0$$

Then homogeneous problem

$$\nabla^2 \omega_H = 0$$

Sep. of variables:

$$w_H(x, y) = X(x) Y(y)$$

$$\Delta^2 w_H = X'' Y + X Y'' = 0$$

$$\frac{X''}{X} = - \frac{Y''}{Y} = \text{const.} = +\lambda^2$$

2 ODEs:

$$Y'' + \lambda^2 Y = 0$$

$$X'' - \lambda^2 X = 0$$

$$Y(y) = A \sin(\lambda y) + B \cos(\lambda y)$$

BC: $w_H|_{y=\pm a/2} = 0$

& symmetric soln. only $\rightarrow A=0$

$$\cos(\lambda \frac{a}{2}) = 0$$



$$\frac{a}{2} \lambda = (2m-1) \frac{\pi}{2}$$

for any $m \in \mathbb{N}$

$$\lambda_m = (2m-1) \frac{\pi}{a}$$

Also:

$$\bar{X}(x) = C e^{\lambda x} + D e^{-\lambda x}$$

Symmetry demands $\bar{X}(x) = \bar{X}(-x)$

$$C = D$$

$$\bar{X}(x) = C (e^{\lambda x} + e^{-\lambda x}) = 2C \cosh \lambda x$$

So the complete soln. is given by

$$w_H \approx \cosh(\lambda_m x) \cos(\lambda_m y)$$

& since the eqn. is linear this holds for arbitrary linear combinations of these facts. for all values of m .

$$w_H(x, y) = \sum_{m=1}^{\infty} A_m \cosh(\lambda_m x) \cos(\lambda_m y)$$

$$\text{for } \lambda_m = (2m-1) \frac{\pi}{a}$$

$$\left. \begin{array}{l} \end{array} \right\} w = w_p + w_H$$

$$w = \frac{1}{2} \frac{G}{\mu} \left(y^2 - \left(\frac{a}{2} \right)^2 \right) + \sum_{m=1}^{\infty} A_m \cosh(\lambda_m x) \cos(\lambda_m y)$$

The B.C. @ $y = \pm \frac{a}{2}$ are already fulfilled thanks to our clever choice of the particular soln. & the choice of λ_m .

The soln. is symmetric in x so we only need to fulfill

$$w(y, x = \frac{a}{2}) = 0$$

This provides an eqn. for the A_m .

$$\frac{1}{2} \frac{G}{\mu} \left(\left(\frac{a}{2} \right)^2 - y^2 \right) = \sum_{m=1}^{\infty} A_m \cosh(\lambda_m \frac{a}{2}) \cos(\lambda_m y)$$

Fourier expand LHS:

$$\frac{1}{2} \frac{G}{\mu} \left(\left(\frac{a}{2} \right)^2 - y^2 \right) = \sum B_m \cos \left((2m-1) \frac{\pi}{a} y \right)$$

$$\int_{-a/2}^{a/2} \cos \left((2m-1) \frac{\pi}{a} y \right) \cos \left((2n-1) \frac{\pi}{a} y \right) dy$$

$$= \begin{cases} \frac{a}{2} & \text{for } m=n \\ 0 & \text{for } m \neq n \end{cases}$$

$$\int_{-a/2}^{a/2} \frac{1}{2} \frac{G}{\mu} \left(\left(\frac{a}{2} \right)^2 - y^2 \right) \cos \left((2m-1) \frac{\pi}{a} y \right) dy$$

(opposite sign to result given)

$$= - \frac{2a^3 G (-1)^m}{(2m-1)^3 \pi^3 \mu} = \frac{a}{2} B_m$$

$$B_m = \frac{4a^2 G (-1)^m}{(2m-1)^3 \pi^3 \mu}$$

$$\sum_1 B_m \cos(\lambda_m y) = \sum_1 A_m \cosh(\lambda_m \frac{a}{2}) \cos(\lambda_m y)$$

compare coeffs:

$$A_m = \frac{B_m}{\cosh(\lambda_m (\frac{a}{2}))}$$

$$\omega = \frac{1}{2} \frac{G}{\mu} \left(y^2 - \left(\frac{a}{2} \right)^2 \right) - \sum_{m=1}^{\infty} \frac{4a^2 G (-1)^m \cosh\left(\frac{(2m-1)\pi}{a} x\right)}{(2m-1)^3 \pi^3 \mu \cosh\left(\frac{(2m-1)\pi}{2}\right)} \times \cos\left(\frac{(2m-1)\pi}{a} y\right)$$