

## The Vorticity equation

- Use of the streamfunction allows us to automatically satisfy the continuity equation. Now we will try to transform the momentum equation

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u}$$

into an equation for the vorticity  $\boldsymbol{\omega} = \nabla \times \mathbf{u}$ .

- Here are a few results that we will use in the derivation:

$$\frac{1}{2} \nabla (\mathbf{u} \cdot \mathbf{u}) = (\mathbf{u} \cdot \nabla) \mathbf{u} + \mathbf{u} \times (\nabla \times \mathbf{u}), \quad (1)$$

$$\nabla \times \nabla \phi = 0, \quad (2)$$

and

$$\nabla \times (\mathbf{u} \times \boldsymbol{\omega}) = (\boldsymbol{\omega} \cdot \nabla) \mathbf{u} - (\mathbf{u} \cdot \nabla) \boldsymbol{\omega} + \mathbf{u} \underbrace{\nabla \cdot \boldsymbol{\omega}}_0 - \boldsymbol{\omega} \underbrace{\nabla \cdot \mathbf{u}}_0, \quad (3)$$

where the last two terms vanish because

$$\nabla \cdot \boldsymbol{\omega} = \nabla \cdot (\nabla \times \mathbf{u}) = \text{div curl } \mathbf{u} = 0$$

and

$$\nabla \cdot \mathbf{u} = 0.$$

- First, we use (1) in the momentum equation to obtain

$$\frac{\partial \mathbf{u}}{\partial t} + \frac{1}{2} \nabla(\mathbf{u} \cdot \mathbf{u}) - \mathbf{u} \times \underbrace{(\nabla \times \mathbf{u})}_{\boldsymbol{\omega}} = -\frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u}$$

i.e.

$$\underbrace{\frac{\partial \mathbf{u}}{\partial t} + \frac{1}{2} \nabla(\mathbf{u} \cdot \mathbf{u}) - \mathbf{u} \times \boldsymbol{\omega}}_{LHS} = \underbrace{-\frac{1}{\rho} \nabla p + \nu \nabla^2 \mathbf{u}}_{RHS}$$

- Now take the curl of the LHS:

$$\nabla \times LHS = \frac{\partial}{\partial t} \underbrace{(\nabla \times \mathbf{u})}_{\boldsymbol{\omega}} + \frac{1}{2} \underbrace{\nabla \times \nabla(\mathbf{u} \cdot \mathbf{u})}_{0 \text{ because of (2)}} - \underbrace{\nabla \times (\mathbf{u} \times \boldsymbol{\omega})}_{\text{see (3)}}$$

i.e.

$$\nabla \times LHS = \frac{\partial \boldsymbol{\omega}}{\partial t} - (\boldsymbol{\omega} \cdot \nabla) \mathbf{u} + (\mathbf{u} \cdot \nabla) \boldsymbol{\omega}$$

- ...and the RHS:

$$\nabla \times RHS = -\frac{1}{\rho} \underbrace{\nabla \times \nabla p}_{0 \text{ because of (2)}} + \nu \nabla^2 \underbrace{(\nabla \times \mathbf{u})}_{\boldsymbol{\omega}}$$

- Now combine the remaining non-zero terms

$$\underbrace{\frac{\partial \boldsymbol{\omega}}{\partial t} + (\mathbf{u} \cdot \nabla) \boldsymbol{\omega}}_{D\boldsymbol{\omega}/Dt} - (\boldsymbol{\omega} \cdot \nabla) \mathbf{u} = \nu \nabla^2 \boldsymbol{\omega}.$$

- The resulting equation is the *vorticity transport equation*

$$\frac{D\boldsymbol{\omega}}{Dt} = (\boldsymbol{\omega} \cdot \nabla)\mathbf{u} + \nu\nabla^2\boldsymbol{\omega} \quad (4)$$

which shows that the rate of change of the vorticity of material particles,  $D\boldsymbol{\omega}/Dt$ , is controlled by ‘vortex stretching’ (described by  $(\boldsymbol{\omega} \cdot \nabla)\mathbf{u}$ ; this is a familiar result from inviscid fluid mechanics) and by diffusion (described by  $\nu\nabla^2\boldsymbol{\omega}$ ). The diffusion of vorticity only occurs in viscous flows.

- For 3D flows, the first term on the RHS in (4) represents vortex stretching: velocity gradients lead to a change in the rate of rotation of material particles.
- Note that for 2D flows, vortex stretching is absent since  $\mathbf{u} = u(x, y)\mathbf{e}_x + v(x, y)\mathbf{e}_y$  and  $\boldsymbol{\omega} = \omega(x, y)\mathbf{e}_z$  and therefore  $(\boldsymbol{\omega} \cdot \nabla)\mathbf{u} = 0$ .
- The vorticity transport equation provides an interesting interpretation of the kinematic viscosity  $\nu$ : The kinematic viscosity is the diffusion coefficient for the diffusion of vorticity.
- Many phenomena in viscous fluid mechanics can be interpreted in terms of the diffusion of vorticity but this is (unfortunately) beyond the scope of this course.

- For 2D flows, the vorticity transport equation

$$\frac{D\omega}{Dt} = \nu \nabla^2 \omega$$

together with the equation for the vorticity in terms of the streamfunction

$$\omega = -\nabla^2 \psi$$

and

$$u = \partial\psi/\partial y \quad \text{and} \quad v = -\partial\psi/\partial x$$

provide the streamfunction-vorticity formulation of the Navier-Stokes equations, which consists of only two PDEs for the scalars  $\omega$  and  $\psi$  rather than the three equations for  $u$ ,  $v$  and  $p$  in the ‘primitive variable’ form.

- Scaling arguments show that in the limit of zero Reynolds number, only one fourth-order PDE for the streamfunction  $\psi$  needs to be solved, namely the biharmonic equation

$$\nabla^4 \psi = 0,$$

where

$$\nabla^4 = \frac{\partial^4}{\partial x^4} + 2\frac{\partial^4}{\partial x^2 \partial y^2} + \frac{\partial^4}{\partial y^4}.$$

- This can also be shown directly by taking the curl of the Stokes equations.