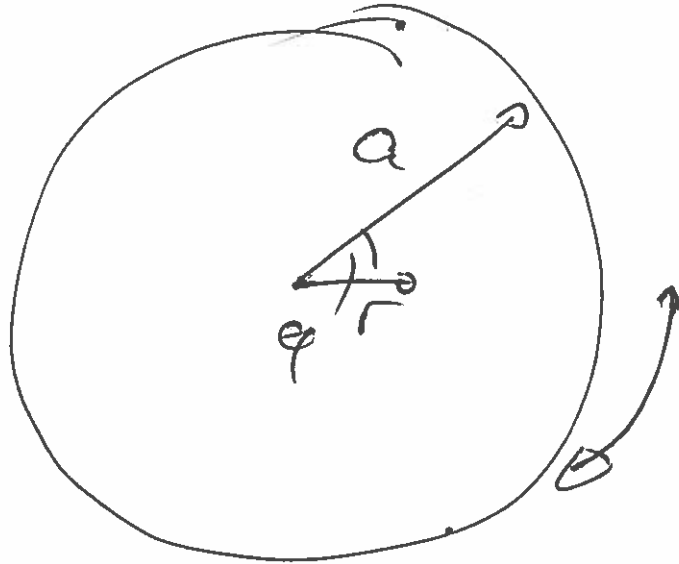


Example:

Oscillating  
Cylinder



$$\underline{y} = \underbrace{\Omega a}_{\omega} \cos(\Omega t) \underline{e}_\varphi$$

$$\underline{y} = U \cos(\Omega t) \underline{e}_\varphi$$

Assumptions:

$$u = u_r = 0$$

$$v = v_\varphi = U(r, t)$$

$$p = p(r, t)$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial r} + \frac{v \partial u}{r \partial \varphi} + w \frac{\partial u}{\partial z} - \frac{v^2}{r} = -\frac{1}{\rho} \frac{\partial P}{\partial r} + \nu \left[ \nabla^2 u - \frac{u}{r^2} - \frac{2}{r^2} \frac{\partial v}{\partial \varphi} \right],$$

$$\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial r} + \frac{v \partial v}{r \partial \varphi} + w \frac{\partial v}{\partial z} + \frac{uv}{r} = -\frac{1}{\rho r} \frac{\partial P}{\partial \varphi} + \nu \left[ \nabla^2 v - \frac{v}{r^2} + \frac{2}{r^2} \frac{\partial u}{\partial \varphi} \right],$$

$$\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial r} + \frac{v \partial w}{r \partial \varphi} + w \frac{\partial w}{\partial z} = -\frac{1}{\rho} \frac{\partial P}{\partial z} + \nu \nabla^2 w,$$

0 = 0

$$\text{div } \mathbf{u} = \frac{1}{r} \frac{\partial}{\partial r} (ru) + \frac{1}{r} \frac{\partial v}{\partial \varphi} + \frac{\partial w}{\partial z} = 0.$$

0 = 0

where

$$\nabla^2 \equiv \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2}.$$

$\varphi$ -mom. eqn:

3

$$\frac{\partial u}{\partial t} = \nu \left( \underbrace{\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r}}_{\nabla^2 u(r,t)} - \frac{u}{r^2} \right) \quad (*)$$

IVP need IC  ~~$\phi$~~   
or  
search for  
time-periodic solen.

azimuthal

$$u(r,t) = \text{Re} \left( \bar{V}(r) e^{i\Omega t} \right)$$

into (\*)

$$i\Omega \bar{V}(r) e^{i\Omega t} = \nu e^{i\Omega t} \left( \bar{V}''(r) + \frac{1}{r} \bar{V}'(r) - \frac{\bar{V}(r)}{r^2} \right)$$

$$\boxed{r^2 \bar{V}'' + r \bar{V}' + \left( -\frac{i\Omega r^2}{\nu} - 1 \right) \bar{V} = 0}$$

compare

$$\xi^2 f''(\xi) + \xi f'(\xi) + (\xi^2 - n^2) f(\xi) = 0$$

(\*\*)

is a Bessel ODE.

(4)

coordinate transformation:

$$n=1$$

$$\xi^2 = -\frac{i\Omega}{\nu} r^2$$

$$r = \sqrt{\frac{i\nu}{\Omega}} \xi$$

Now note that derivs in ~~(\*\*)~~ are of the form  $r^n \frac{df}{dr^n}$ .

Such expressions are invariant under re-scalings,  $r \rightarrow \alpha r$

$$\begin{array}{c} r \rightarrow \alpha r \\ \uparrow \\ \text{Constant} \end{array}$$

$$r = \sqrt{\frac{i\nu}{\Omega}} \xi$$

$$\frac{dr}{d\xi} = \alpha$$

$$\frac{df}{dr} = \frac{1}{\frac{dr}{d\xi}}$$

$$\frac{df}{d\xi} = \alpha \frac{df}{dr}$$

$$\boxed{r \frac{df}{dr} = \alpha \xi \frac{1}{\alpha} \frac{df}{d\xi} = \xi \frac{df}{d\xi}}$$

$$r^2 \frac{d^2 h}{dr^2} = \alpha^2 \cancel{r}^2 \frac{d}{d\cancel{r}} \left( \frac{dh}{d\cancel{r}} \right)$$

$$= \cancel{\alpha^2} \cancel{r}^2 \frac{1}{\cancel{r}} \frac{d}{d\cancel{r}} \left( \frac{1}{\cancel{r}} \frac{dh}{d\cancel{r}} \right)$$

$$r^2 \frac{d^2 h}{dr^2} = \cancel{r}^2 \frac{d^2 h}{d\cancel{r}^2}$$

So with  $r = \sqrt{\frac{i\omega}{\rho}} \xi$  (AA)  
 becomes:

$$\xi^2 V''(\xi) + \xi V'(\xi) + (\xi^2 - 1) V(\xi) = 0$$

Bessel ODE with  $n=1$ .  
 Gen. form. of this ODE is

(lin., 2nd order, homogeneous)

$$f(\xi) = A J_1(\xi) + B Y_1(\xi)$$

Bessel fns of the 1st & 2nd kind of order 1

So gen. soln of (\*\*):

(6)

$$V(r) = A J_1\left(\sqrt{\frac{\rho}{i\omega}} r\right) + B Y_1\left(\sqrt{\frac{\rho}{i\omega}} r\right)$$

BC:  $V(r=0) = U$

$V$  finite as  $r \rightarrow 0$ .

Note:  $Y_1(x) \rightarrow \infty$  as  $x \rightarrow 0$

$\Rightarrow B = 0$

$$V(r) = \frac{U}{J_1\left(\sqrt{\frac{\rho}{i\omega}} a\right)} J_1\left(\sqrt{\frac{\rho}{i\omega}} r\right)$$

$$u(r,t) = \operatorname{Re} \left( \frac{U J_1\left(\sqrt{\frac{\rho}{i\omega}} r\right)}{J_1\left(\sqrt{\frac{\rho}{i\omega}} a\right)} e^{i\omega t} \right)$$

$$\frac{u(r,t)}{U} = \operatorname{Re} \left( \frac{J_1\left(\sqrt{\frac{\rho a^2}{i\omega}} \left(\frac{r}{a}\right)\right)}{J_1\left(\sqrt{\frac{\rho a^2}{i\omega}}\right)} e^{i\omega t} \right)$$

Note: The veloc. profile only depends on the non-dimensional parameter

$$\frac{\Omega a^2}{\nu} \sim \frac{\text{inertial}}{\text{viscosity}}$$

(See animation on webpage)

Consider the case  $\Omega \rightarrow 0$

Fact:

$$J_1(x) = \frac{1}{2}x - \frac{1}{16}x^3 + \dots$$

As  $\Omega \rightarrow 0$

$$u(r,t) = \text{Re} \left( \frac{1}{2} \sqrt{\frac{\Omega}{i\nu}} r + \dots + \frac{1}{2} \sqrt{\frac{\Omega}{i\nu}} a + \dots \right) e^{i\Omega t}$$

$$u(r,t) \rightarrow \frac{\Gamma}{a} u \cos(\Omega t)$$

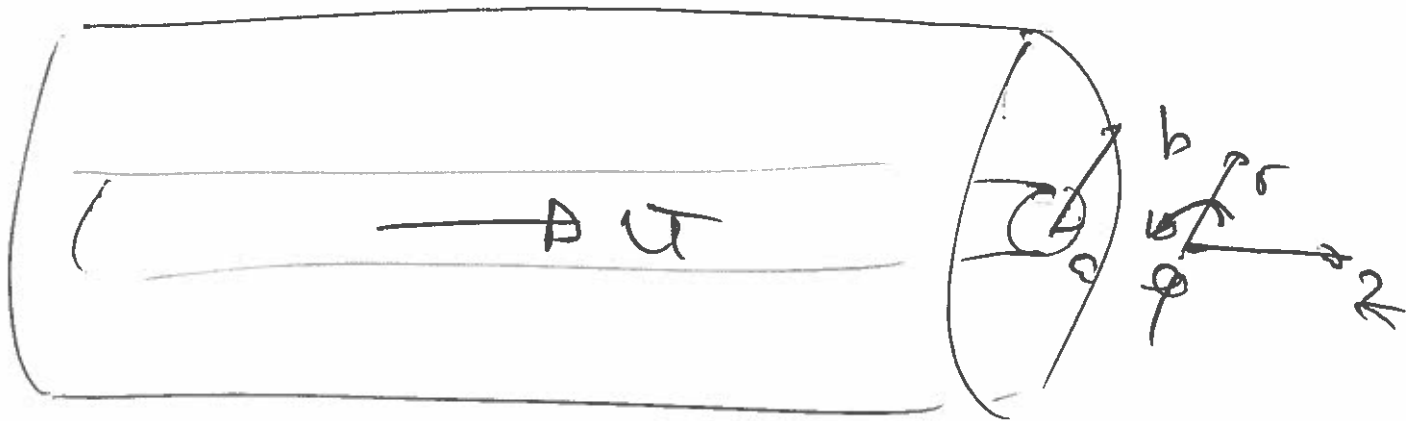
⇒ rigid body rotation!  $\mathcal{L}$

quasi-steady

(as expected)



X-class

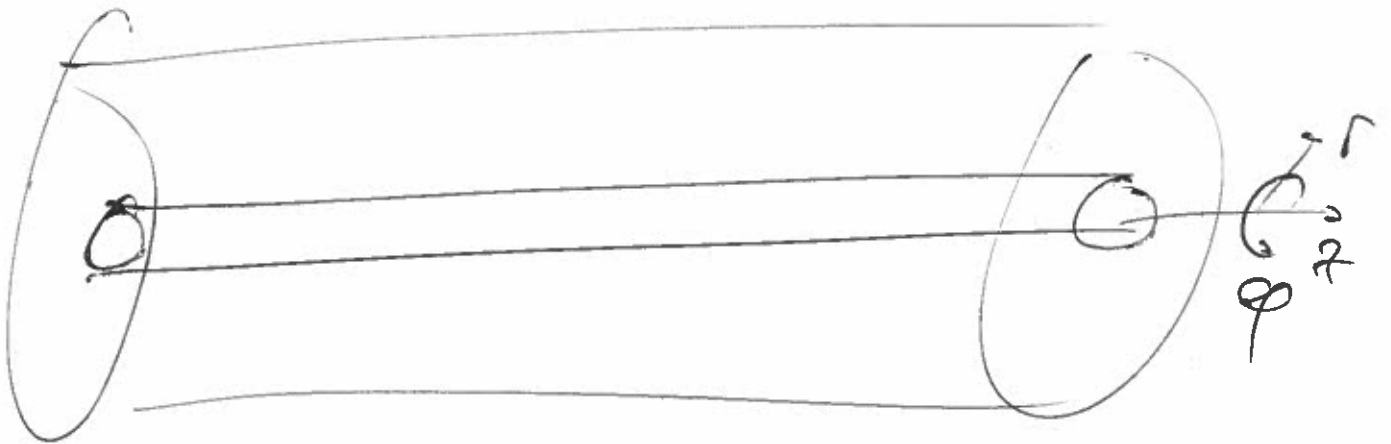


$$\nabla p = G \underline{e}_z$$

$$\underline{u}(r, z, \phi, t) = \omega(r) \underline{e}_\phi$$

$$t_1 = \tau_{r\theta} n_\theta$$

$$t_2 = ?$$



Traction by rod on fluid?

at  $r = 0$ :  $\underline{n} = -\underline{e}_r$

$$\underline{n} = n_r \underline{e}_r + n_\theta \underline{e}_\theta + n_z \underline{e}_z$$

$$n_r = -1 \quad n_\theta = n_z = 0$$

$$\tau_{ij} = -p \delta_{ij} + 2\mu \underline{\underline{\epsilon}}_{ij}$$

$$\underline{u} = \omega(r) \underline{e}_z$$