

EXAMPLE SHEET II

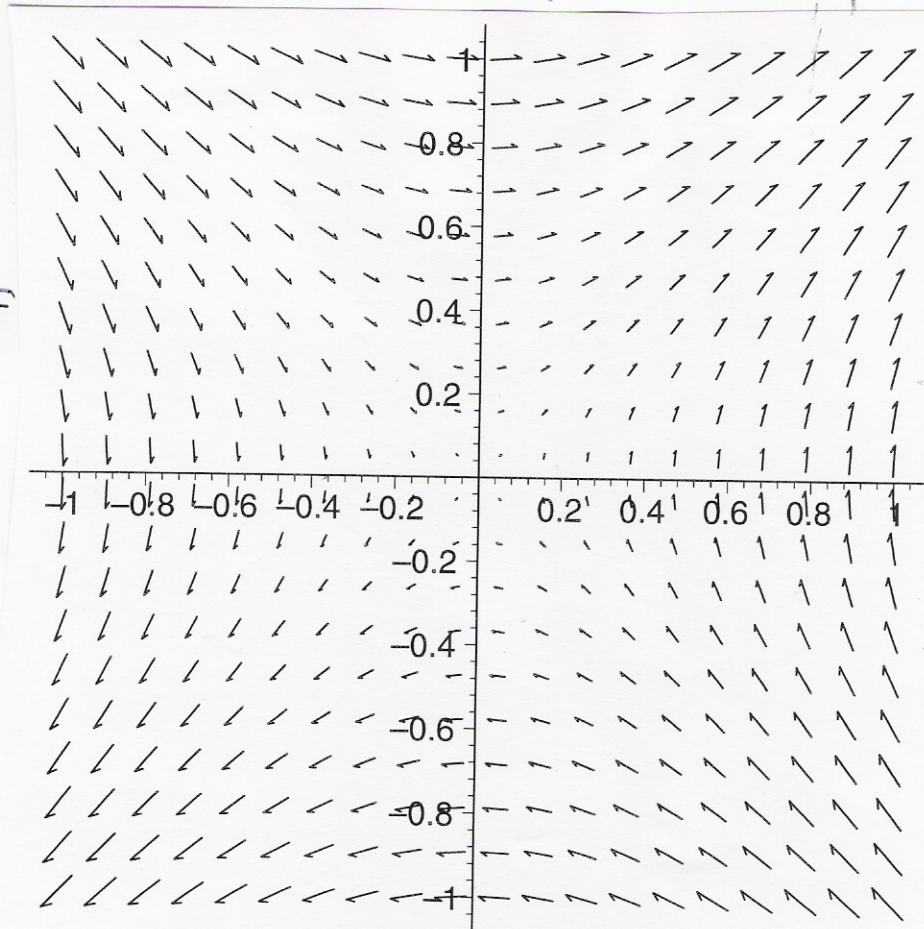
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1) a)

Second
diagonal:

$$x_1 = -x_2 = \omega$$

$$u = a \omega \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$



$x_2 = 0:$
 $u_1 = 0$
 $u_2 = a x_1$

main
Diagonal:

$$x_1 = x_2 = \omega$$

$$u = a \omega \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$x_1 = 0:$ $u_1 = a x_2, u_2 = 0$

b) Eqn. of motion:

$$u_1 = a x_2^p = \frac{\partial x_1^p}{\partial t} \quad (1)$$

$$u_2 = a x_1^p = \frac{\partial x_2^p}{\partial t} \quad (2)$$

$\frac{\partial}{\partial t}$ (1) into (2)

$$a x_1^p = \frac{\partial x_2^p}{\partial t} = \frac{1}{a} \frac{\partial^2 x_1^p}{\partial t^2}$$

$$\ddot{x}_i^p - a^2 x_i^p = 0$$

$$x_i^p(t) = A e^{at} + B e^{-at}$$

Also from (1):

$$x_2^p = \frac{1}{a} \frac{\partial x_i^p}{\partial t} = A e^{at} - B e^{-at}$$

initial cond:

$$\text{At time } t=0: x_i^p = \bar{x}_i$$

$$\left. \begin{aligned} A+B &= \bar{x}_1 \\ A-B &= \bar{x}_2 \end{aligned} \right\} \begin{aligned} A &= \frac{1}{2}(\bar{x}_1 + \bar{x}_2) \\ B &= \frac{1}{2}(\bar{x}_1 - \bar{x}_2) \end{aligned}$$

$$\begin{aligned} x_1^p(t) &= \frac{1}{2}(\bar{x}_1 + \bar{x}_2) e^{at} + \frac{1}{2}(\bar{x}_1 - \bar{x}_2) e^{-at} \\ x_2^p(t) &= \frac{1}{2}(\bar{x}_1 + \bar{x}_2) e^{at} - \frac{1}{2}(\bar{x}_1 - \bar{x}_2) e^{-at} \end{aligned}$$

This gives the spatial posn. of a particle in the flow at time t which was at \bar{x}_i at time $t=0$.

c) The acceleration of the individual particle is

$$a_i = \frac{\partial^2 x_i^p}{\partial t^2}$$

w.l.o.g. we can consider the situation at $t=0$ (the flow is steady!).

The particle is at $x_i = \bar{x}_i$ & we have

$$a_1 = \frac{\partial^2 x_1^p}{\partial t^2} (x_i = \bar{x}_i) = a^2 \bar{x}_1 = a^2 x_1$$

$$a_2 = \frac{\partial^2 x_2^p}{\partial t^2} (x_i = \bar{x}_i) = a^2 \bar{x}_2 = a^2 x_2$$

d) using the substantial derivative, we obtain:

$$a_i = \frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j}$$

Steady flow

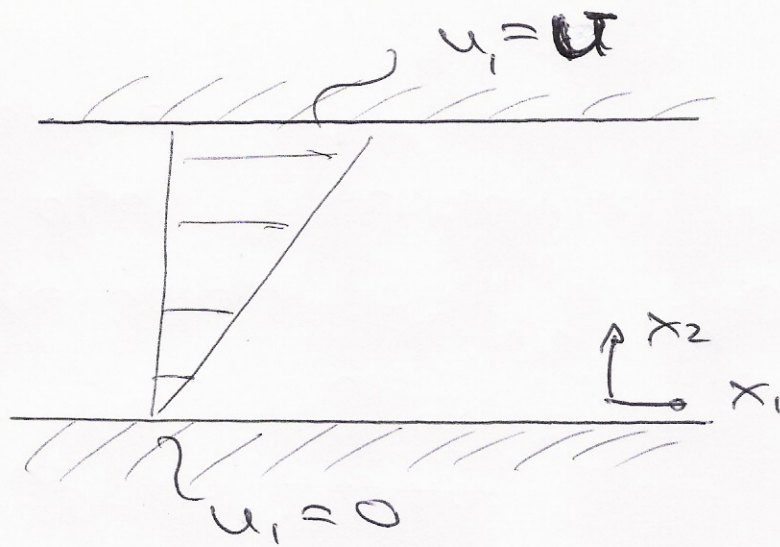
$$a_1 = u_1 \frac{\partial u_1}{\partial x_1} + u_2 \frac{\partial u_1}{\partial x_2} = a x_2 \cdot 0 + a x_1 \cdot a = a^2 x_1 \quad \checkmark$$

$$a_2 = u_1 \frac{\partial u_2}{\partial x_1} + u_2 \frac{\partial u_2}{\partial x_2} = a x_2 \cdot a + a x_1 \cdot 0 = a^2 x_2 \quad \checkmark$$

(4)

Remember: in both cases, we are interested in the acceleration of the particle that is currently at a given spatial position. This is not the same as $\frac{du}{dt}$ (which is zero in the present case!)

2) a)



$$u_1 = U x_2$$

$$u_2 = 0$$

This flow field could be the flow in a 2D channel with a stationary bottom wall (at $x_2=0$) & a moving top wall which moves to the right with velocity U .

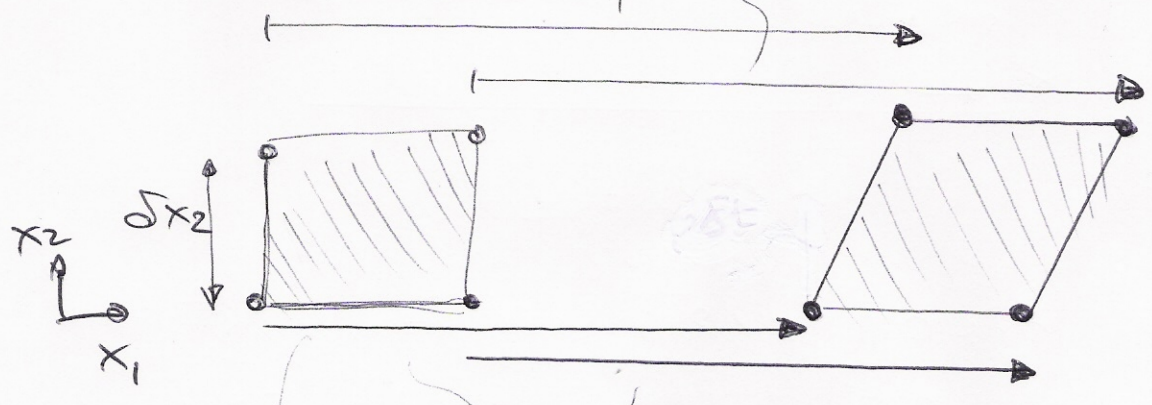
$$b) \quad \varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

$$\omega_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right)$$

$$\frac{\partial u_i}{\partial x_j} = \begin{array}{c|cc} & j=1 & j=2 \\ \hline i=1 & 0 & U \\ i=2 & 0 & 0 \end{array}$$

$$\varepsilon_{ij} = \begin{pmatrix} 0 & \frac{1}{2}U \\ \frac{1}{2}U & 0 \end{pmatrix} \quad \omega_{ij} = \begin{pmatrix} 0 & \frac{1}{2}U \\ -\frac{1}{2}U & 0 \end{pmatrix}$$

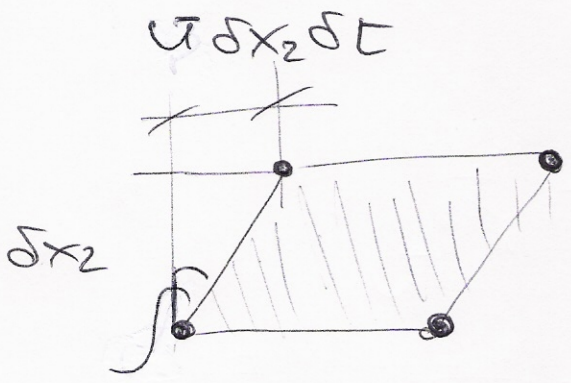
$$u_1(x_2 + \delta x_2) \delta t = u(x_2 + \delta x_2) \delta t$$



NO
AXIAL
EXTENSION
($\epsilon_{11} = 0$)

$$u_1(x_2 = \bar{x}_2) \delta t = u \bar{x}_2$$

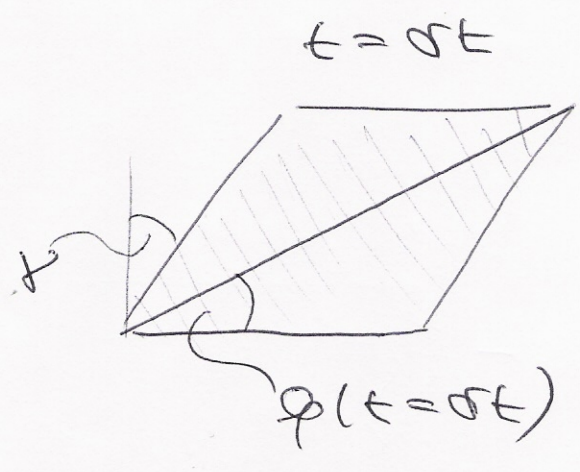
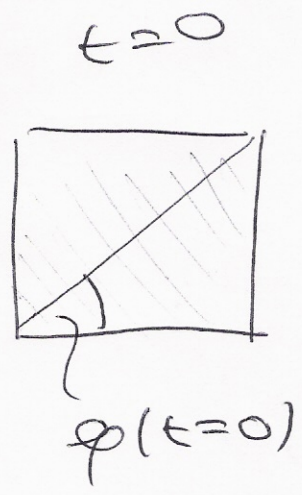
(lower edge at $x_2 = \bar{x}_2$, say)



$$\text{top } \mu \approx \mu = \frac{u \delta x_2 \delta t}{\delta x_2}$$

$$\frac{D\mu}{Dt} = u = 2\epsilon_{12} \quad \checkmark \quad \text{just as claimed } \text{☺}$$

Also consider the rotation of a (square) element by following its main diagonal:



($\varphi = \frac{\pi}{4}$ for a square elem.)

$$\varphi = \frac{1}{2} \left(\frac{\pi}{2} - \theta \right)$$

$$\frac{D\varphi}{Dt} = -\frac{1}{2} \theta$$

The main diagonal rotates about the x_3 axis (which comes out of the paper!) with a (neg.) angular velocity

$$\frac{D\varphi}{Dt} = \Omega_3 = -\frac{1}{2} \theta = \omega_{21}$$

✓ just as claimed in the lecture

