

$$u = \frac{\partial \psi}{\partial y}$$

$$v = -\frac{\partial \psi}{\partial x}$$

(1)

$$\frac{D\omega}{Dt} = (\underbrace{\omega \cdot \nabla}_{\text{0 in 2D}}) \underline{u} + \nu \nabla^2 \underline{\omega}$$

0 in 2D; also

$$\underline{\omega} = \omega \underline{e}_z$$

Scale:

Length: a ; velocity-scale: U

$$\underline{u} = U \underline{u}^*$$

$$\underline{r} = a \underline{r}^*$$

$$t = \frac{a}{U} t^*$$

$$\underline{\omega} = \nabla \times \underline{u} = \frac{U}{a} \underline{\omega}^*$$

$$\frac{U^2}{a^2} \frac{\partial \tilde{\omega}_i}{\partial \tilde{t}} + \frac{U^2}{a^2} \tilde{u}_j \frac{\partial \tilde{\omega}_i}{\partial \tilde{x}_j} = \frac{U^2}{a^2} \tilde{\omega}_i \frac{\partial \tilde{u}_i}{\partial \tilde{x}_j} +$$

$$\frac{U}{a^3} \nu \frac{\partial^2 \tilde{\omega}_i}{\partial \tilde{x}_j^2}$$

$$\frac{Ua}{\nu} \frac{D\tilde{\omega}}{D\tilde{t}} = \frac{Ua}{\nu} (\tilde{\omega} \cdot \tilde{\nabla}) \tilde{u} + \tilde{\nabla}^2 \tilde{\omega}$$

As $Re \rightarrow 0$:

(2)

$$\nabla^2 \underline{\omega} = \underline{0}; \quad \nabla^2 \underline{\omega} = \underline{0}$$

Aside: Start with Stokes eqns:

$$\nabla p = \mu \nabla^2 \underline{u}$$

Take curl:

$$\underbrace{\nabla \times \nabla p}_{\underline{0}} = \mu \nabla^2 \underbrace{\nabla \times \underline{u}}_{\underline{\omega}}$$

$$\nabla^2 \underline{\omega} = \underline{0}$$

Now consider 2D flow:

$$\underline{\omega} = \omega \underline{e}_z$$

$$\nabla^2 \omega = 0 \quad \text{also} \quad \omega = -\nabla^2 \psi$$

2D Stokes flow

\Rightarrow

$$\boxed{\nabla^2 \nabla^2 \psi = 0}$$

$$\nabla^4 \psi = 0$$

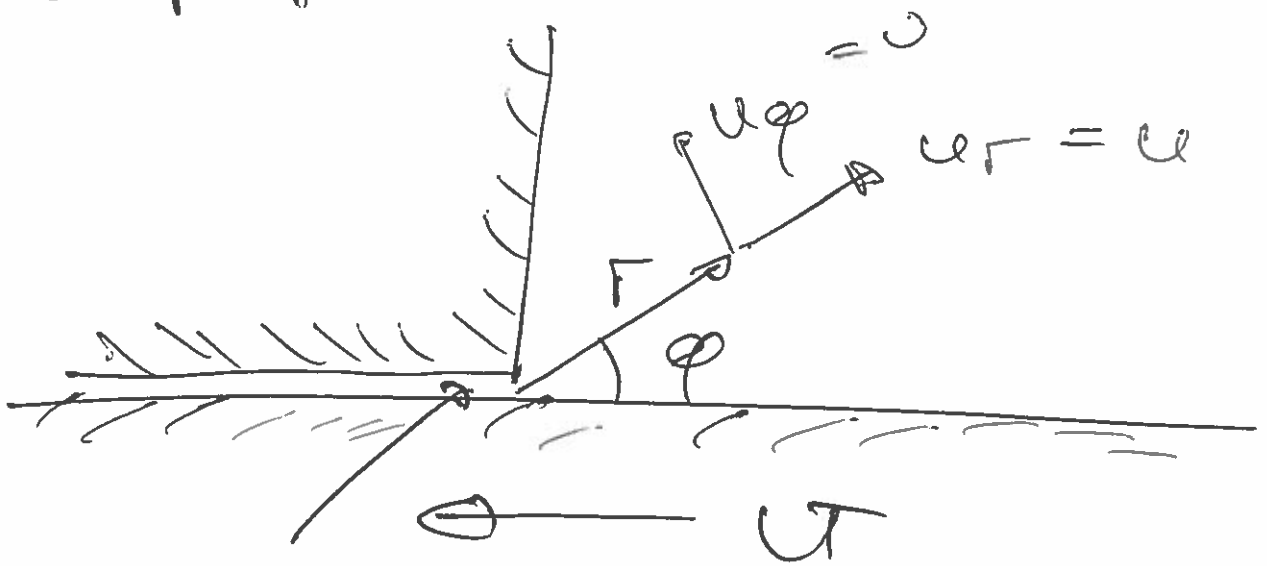
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$$\nabla^4 = \frac{\partial^4}{\partial x^4} + 2 \frac{\partial^4}{\partial x^2 \partial y^2} + \frac{\partial^4}{\partial y^4}$$

one scalar ψ 'th order linear PDE for $\psi(x, y)$

Example:

Scraping flow



zero
gap width.

Assume: slow, viscous flow

\Rightarrow Stokes eqns are valid since $Re \ll 1$

$$\nabla^4 \psi = 0$$

Use polar coords.

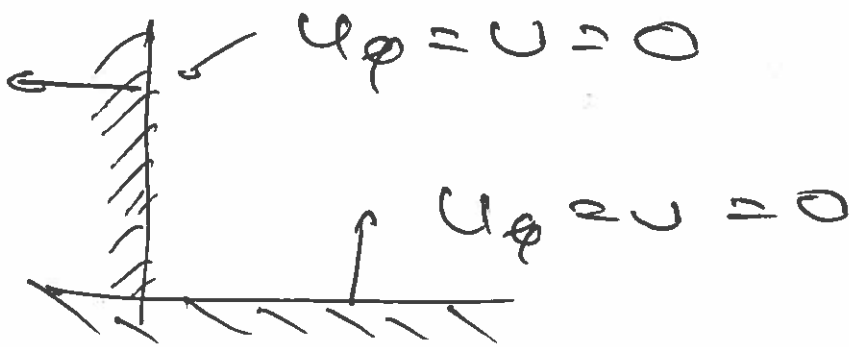
(4)

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2}$$

$$u = u_r = \frac{1}{r} \frac{\partial \psi}{\partial \phi}$$

$$v = u_\phi = - \frac{\partial \psi}{\partial r}$$

BC: Impermeability:



So:

$$v = - \frac{\partial \psi}{\partial r} = 0 \quad \text{at } \phi = 0 \text{ to } \pi$$

$\psi(\phi = 0) = \text{const} = C_1$

Also:

$$v = - \frac{\partial \psi}{\partial r} = 0 \quad \text{at } \phi = \frac{\pi}{2} \text{ to } \frac{3\pi}{2}$$

$\psi(\phi = \frac{\pi}{2}) = \text{const} = C_2$

Note: zero flux into

(5)

corner $\Rightarrow \psi$ continuous

$$\Rightarrow \sigma_1 = \sigma_2 = \sigma = 0$$

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w.l.o.f.
for u

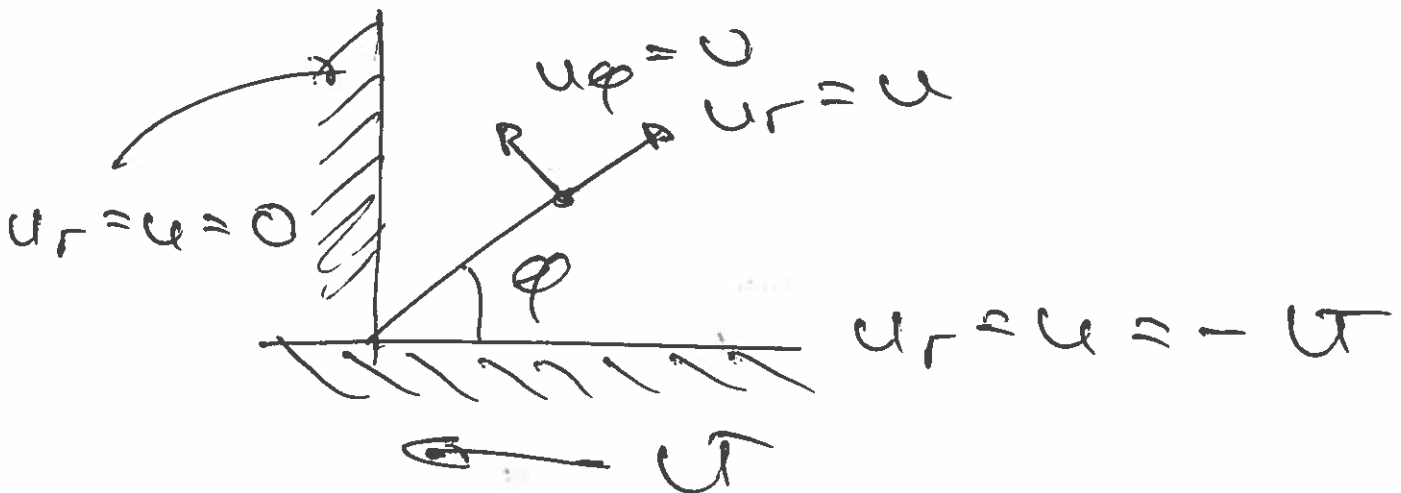
$$\psi(\varphi=0) = 0$$

(1)

$$\psi(\varphi = \frac{\pi}{2}) = 0$$

(2)

No slip:



$$u = 0 = \frac{1}{r} \frac{d\psi}{d\varphi} \quad \text{at } \varphi = \frac{\pi}{2} \quad (3)$$

$$u = -U = \frac{1}{r} \frac{d\psi}{d\varphi} \quad \text{at } \varphi = 0 \quad (4)$$

want $\psi(r, \varphi)$

(6)

Try: $\psi(r, \varphi) = g(r) f(\varphi)$

Now BC (3) $f(\varphi)$ show

that $\frac{1}{r} \frac{\partial \psi}{\partial \varphi}$ should be

indep. of r . (at least of $\varphi=0$,
 $\varphi = \frac{\pi}{2}$)

This suggest to try

$$g(r) = Ur$$

Ansatz:

$$\psi(r, \varphi) = Ur f(\varphi)$$

BC for $f(\varphi)$:

$$\varphi = 0: \psi = 0$$

$$\varphi = \frac{\pi}{2}: \psi = 0$$

$$\varphi = 0: \frac{1}{r} \frac{\partial \psi}{\partial \varphi} = -U$$

$$\varphi = \frac{\pi}{2}: \frac{1}{r} \frac{\partial \psi}{\partial \varphi} = 0$$

$$f(0) = 0$$

$$f\left(\frac{\pi}{2}\right) = 0$$

$$f'(0) = -1$$

$$f'\left(\frac{\pi}{2}\right) = 0$$

$$\text{PDE: } \nabla^2 \nabla^2 \psi = 0 \quad (7)$$

$$\nabla^2 \psi = \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \right) \psi r f(\phi)$$

$$= \frac{\psi}{r} f + \frac{\psi}{r} f'' = \psi r^{-1} (f + f'')$$

$$\nabla^2 \nabla^2 \psi = \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \right) \psi r^{-1} (f + f'')$$

$$= \psi \left\{ (-1)(-2) r^{-3} (f + f'') + (-1) r^{-3} (f + f'') \right. \\ \left. + r^{-3} (f'' + f''') \right\}$$

$$= \frac{\psi}{r^3} \left\{ (2-1) f + (2-1+1) f'' + f''' \right\}$$

$$\nabla^4 \psi = \frac{\psi}{r^3} (f''' + 2f'' + f) = 0$$

4th order, linear, homof, const. coeff, r, ϕ

$$f(\varphi) \sim e^{\lambda \varphi}$$

(8)

char. poly:

$$\lambda^4 + 2\lambda^2 + 1 = 0$$

$$(\lambda^2 + 1)^2 = 0$$

$\lambda_{1234} = \pm i$ repeated roots

$$f(\varphi) = ~~f(\varphi)~~$$

$$A \sin \varphi + B \cos \varphi +$$

$$C \varphi \sin \varphi + D \varphi \cos \varphi$$

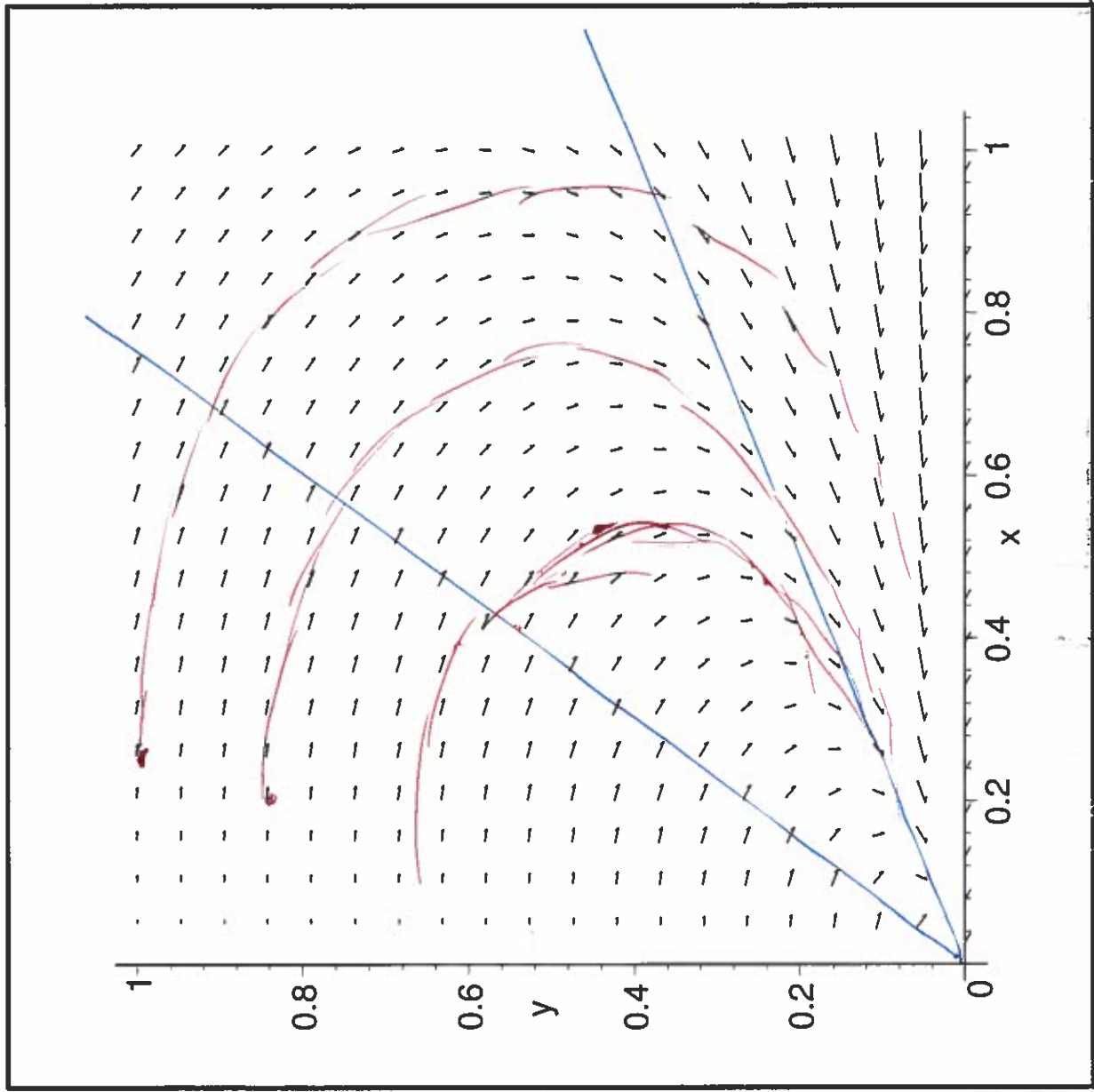
1st year!!!!

A, B, C, D from BC...

$$\psi = \frac{U \Gamma}{\left(\frac{\Gamma}{2}\right)^2 - 1} \left(-\left(\frac{\Gamma}{2}\right)^2 \sin \varphi + \varphi \cos \varphi + \right. \\ \left. + \frac{\Gamma}{2} \varphi \sin \varphi \right)$$

Velocity field for scraping flow at zero Reynolds number:

The vertical wall at $x=0$ is stationary. The horizontal wall at $y=0$ moves to the left with unit velocity.



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So: $u_\phi = u = -\frac{d\psi}{dr}$ } is indep. of r (10)

$u_r = u = \frac{1}{r} \frac{d\psi}{d\phi}$ }

Discussion:

I Nonuniformity of the soln

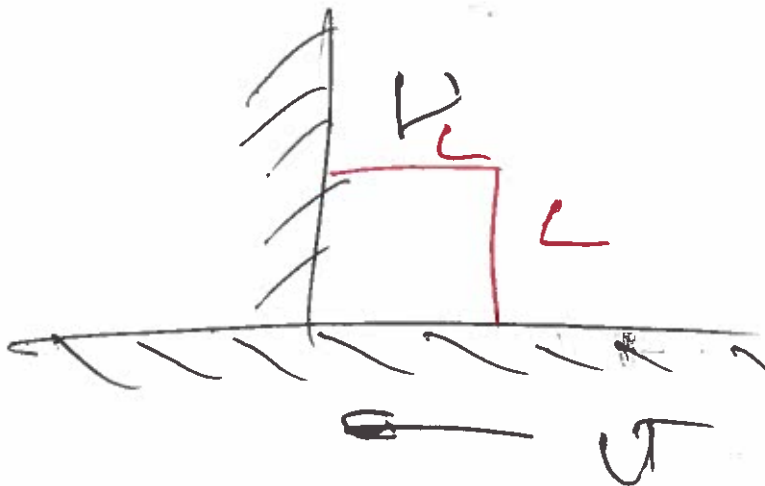
We have assumed $Re \ll 1$

$$Re = \frac{Ua}{\nu}$$

U → veloc of bottom plate

ν → viscosity of fluid

what is a ?



There is no length scale



choose e.g. distance from corner

Choose length scale $L = a$

If a is suff. small then
 Re can become arbitrarily
small BUT equally
at large distances Re
becomes arbitrarily large
 \Rightarrow Stokes eqns invalid.

\Rightarrow & form is only
valid locally.

Need: $Re = \frac{Ua}{\nu} \ll 1$

\Rightarrow $a \ll \frac{\nu}{U}$