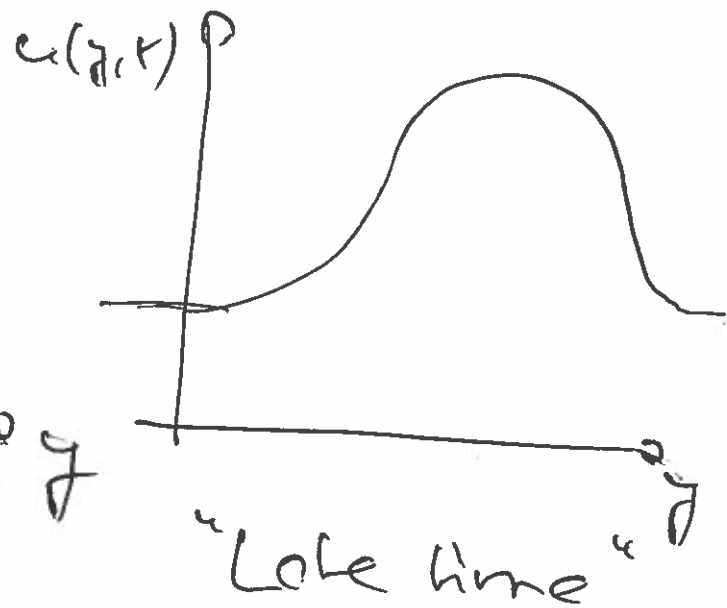


§ Similarity solutions (1)

Often solutions are self-similar; i.e. they have the same "shape" but possibly different scales at different times.

E.g. $u(y, t)$



Such solutions have the form

$$u(y, t) = \underbrace{a(t)}_{\text{amplitude}} \underbrace{f\left(\frac{y}{b(t)}\right)}_{\text{shape}}$$

Sim. solns. don't always exist - try it!

They often reduce PDEs into ODE for $f(\eta)$

where $\eta = \frac{y}{b(x)}$ ← similarity variable.

~~Thus~~ The existence of such solutions is often suggested by dimensional arguments.

Example:

Rayleigh's jerked plate.



Assume: parallel flow (3)

$$\underline{u} = u(y, t) \underline{e}_x ; \nabla p = 0$$

$$\frac{du}{dt} = \nu \frac{d^2 u}{dy^2} \quad (\text{as before})$$

$$\underline{IC}: \quad t = 0: \quad u = 0$$

$$\underline{BC}: \quad \left. \begin{array}{l} u(y=0, t) = U \quad (*) \\ u \rightarrow 0 \text{ as } y \rightarrow \infty \end{array} \right\} \begin{array}{l} \text{For} \\ t \rightarrow 0 \end{array}$$

- Note:
- $u = u(y, t; U, \nu)$
 - PDE is linear & homog.
 - BC/IC are linear in u & homogeneous apart from (*)

$\Rightarrow u(y, t; \nu, U)$ must be a linear fun. of U .
(assuming uniqueness)

Proof: Assume have Poisson \square
for u

$$2 \frac{\partial u}{\partial t} = 2 \nu \frac{\partial^2 u}{\partial y^2}$$

$$t=0: \quad 2u = 0$$

$$\left. \begin{array}{l} 2u(y=0, t) = 2U \\ 2u \rightarrow 0 \quad \text{as } y \rightarrow \infty \end{array} \right\} t > 0$$

\square

$$u(y, t; U, \nu) = U f(y, t; \nu)$$

Now: Dimensions must match! Both sides must have dimensions $\frac{m}{sec}$.

$\Rightarrow f$ must be a non-dim. fct of its arguments.

$$f(y, t; \nu):$$

$$[y] = m$$

$$[t] = \text{sec}$$

$$[\nu] = \frac{m^2}{\text{sec}}$$

$$[\nu] = \frac{\left[\frac{\partial^2 y}{\partial t^2} \right]}{\left[\frac{\partial^2 y}{\partial y^2} \right]} = \frac{\left(\frac{m}{\text{sec}^2} \right)}{\left(\frac{m}{\text{sec}^2 m^2} \right)}$$

$$[\nu] = \frac{m^2}{\text{sec}}$$

$$\gamma = \frac{y^2}{L^2 t}$$

$$\text{or } \frac{L^2 t}{y^2} \dots$$

we choose

$$\gamma = \frac{y}{\sqrt{L t}}$$

So the soln. must
have the form

(6)

$$u(\gamma, t; \nu, \alpha) = \alpha f(\gamma)$$

where $\boxed{\gamma = \gamma (\nu t)^{-1/2}}$

In b PDE:

$$\frac{\partial u}{\partial t} = \alpha \frac{df}{d\gamma} \frac{\partial \gamma}{\partial t}$$

~~where~~

$$\boxed{\frac{\partial u}{\partial t} = \alpha f' \frac{\gamma}{\sqrt{\nu t}} \left(-\frac{1}{2}\right) t^{-3/2}}$$

$$\frac{\partial u}{\partial \gamma} = \alpha f' \frac{\partial \gamma}{\partial \gamma} = \alpha f' (\nu t)^{-1/2}$$

$$\boxed{\frac{\partial^2 u}{\partial \gamma^2} = \alpha f'' \frac{1}{\nu t}}$$

into PDE:

(7)

$$\cancel{u} f' \frac{y}{\sqrt{kt}} \left(\frac{1}{2}\right) t^{-3/2} = \cancel{kt} \cancel{u} f'' \frac{1}{\sqrt{kt}}$$

$\frac{du}{df}$ $\frac{du}{df}$

$$-\frac{1}{2} \frac{y}{\sqrt{kt}} f' = f''$$

η

$$f'' + \frac{1}{2} \eta f' = 0$$

ODE for $f(\eta)$.

$$u(y,t) = u f(\eta)$$
$$\eta = \frac{y}{\sqrt{kt}}$$

BC: $u = uT$ at $y=0$

$$f(\eta=0) = 1$$

$$u \rightarrow 0 \text{ as } y \rightarrow \infty$$

$$f \rightarrow 0 \text{ as } \eta \rightarrow \infty$$

DC: $u \rightarrow \Delta 0$ as $t \rightarrow \Delta 0$

$f \rightarrow \Delta 0$ as $\gamma \rightarrow \Delta 0$

again, phew!