

$$\frac{\partial \hat{u}}{\partial t} = \nu \frac{\partial^2 \hat{u}}{\partial y^2}$$

$$\hat{u}(y=0, t) = U \quad \text{for } t > 0$$

$$\hat{u} \rightarrow 0 \quad \text{as } y \rightarrow \infty$$

$$\hat{u}(y, t=0) = 0$$

$$\hat{u} = \hat{u}(y, t; \nu, U)$$

~~$$\frac{\partial \hat{u}}{\partial t} = \nu \frac{\partial^2 \hat{u}}{\partial y^2}$$~~

~~$$\hat{u}(y=0, t) = U \quad \text{for } t > 0$$~~

~~$$\hat{u} \rightarrow 0 \quad \text{as } y \rightarrow \infty$$~~

~~$$\hat{u}(y, t=0) = 0$$~~

$$u(y, t; \nu, U) = U \hat{u}(y, t; \nu, U)$$

$$u(y, t; \nu, U) = U \underbrace{f(y, t; \nu)}_{\text{dim less.}}$$

dim less.

(2)

combine γ, t & L
into

$$\eta = \frac{\gamma}{\sqrt{L t}}$$

This is the only nondim.
combination of γ, t & L
that is linear in γ .

$$\Rightarrow u(\gamma, t; L, \nu) = \nu f(\eta)$$

into PDE & BC & IC

$$u = \nu f(\underbrace{\gamma (L t)^{-1/2}}_{\eta})$$

$$\frac{\partial u}{\partial t} = \nu \frac{df}{d\eta} \frac{\partial \eta}{\partial t}$$

$$= \nu f' \cdot \frac{\gamma}{\sqrt{L}} \left(-\frac{1}{2}\right) t^{-3/2}$$

$$\frac{\partial u}{\partial \gamma} = \nu \frac{df}{d\eta} \frac{\partial \eta}{\partial \gamma} = \nu f' \frac{1}{\sqrt{L t}} t^{-1/2}$$

$$\frac{\partial^2 u}{\partial y^2} = \text{CT } f'' (\text{LT})^{-1}$$

3

into PDE:

$$\frac{\partial u}{\partial t} = \nu \frac{\partial^2 u}{\partial y^2}$$

$$-\frac{1}{2} \cancel{\text{CT}} \underbrace{\frac{y}{\sqrt{4\nu t}}}_{\eta} \cancel{\text{LT}} f' = \cancel{\nu} \cancel{\text{CT}} \frac{f''}{\cancel{\text{LT}}}$$

3

$$f'' + \frac{1}{2} \eta f' = 0$$

$$\eta = \frac{y}{\sqrt{4\nu t}}$$
$$u = \text{CT } f(\eta)$$

2nd order ODE for $f(\eta)$

BC:

$$u = U_T \text{ at } y = 0$$

$$f = 1 \text{ for } \eta = 0$$

$$u \rightarrow 0 \text{ as } y \rightarrow \infty$$

$$f \rightarrow 0 \text{ as } \eta \rightarrow \infty$$

(*)

IC:

(4)

$$u = 0 \quad \text{as } t \rightarrow 0$$

$$f = 0 \quad \text{as } z \rightarrow \infty$$

same as 2nd BC (*)



Solve:

$$f'' + \frac{1}{2} z f' = 0$$

$$F = f'$$

$$F' + \frac{1}{2} z F = 0$$

$$\frac{F'}{F} = -\frac{1}{2} z$$

$$\ln\left(\frac{F}{F_0}\right) = -\frac{1}{4} z^2$$

↑ const.

$$f' = f = F_0 \exp(-\frac{1}{4}z^2)$$

$$f(z) = A + F_0 \int_{\infty}^z \exp(-\frac{1}{4}\xi^2) d\xi$$

What about lower limit?

Arbitrary, here choose: ∞

$$f(z) = A + B \int_{\infty}^z \exp(-\frac{1}{4}\xi^2) d\xi$$

$-F_0$

Be: ~~$f \rightarrow 0$~~ $f \rightarrow 0$ as $z \rightarrow \infty$.

$$\Rightarrow A = 0$$

$$f(z=0) = 1 = B \int_0^{\infty} \exp(-\frac{1}{4}\xi^2) d\xi$$

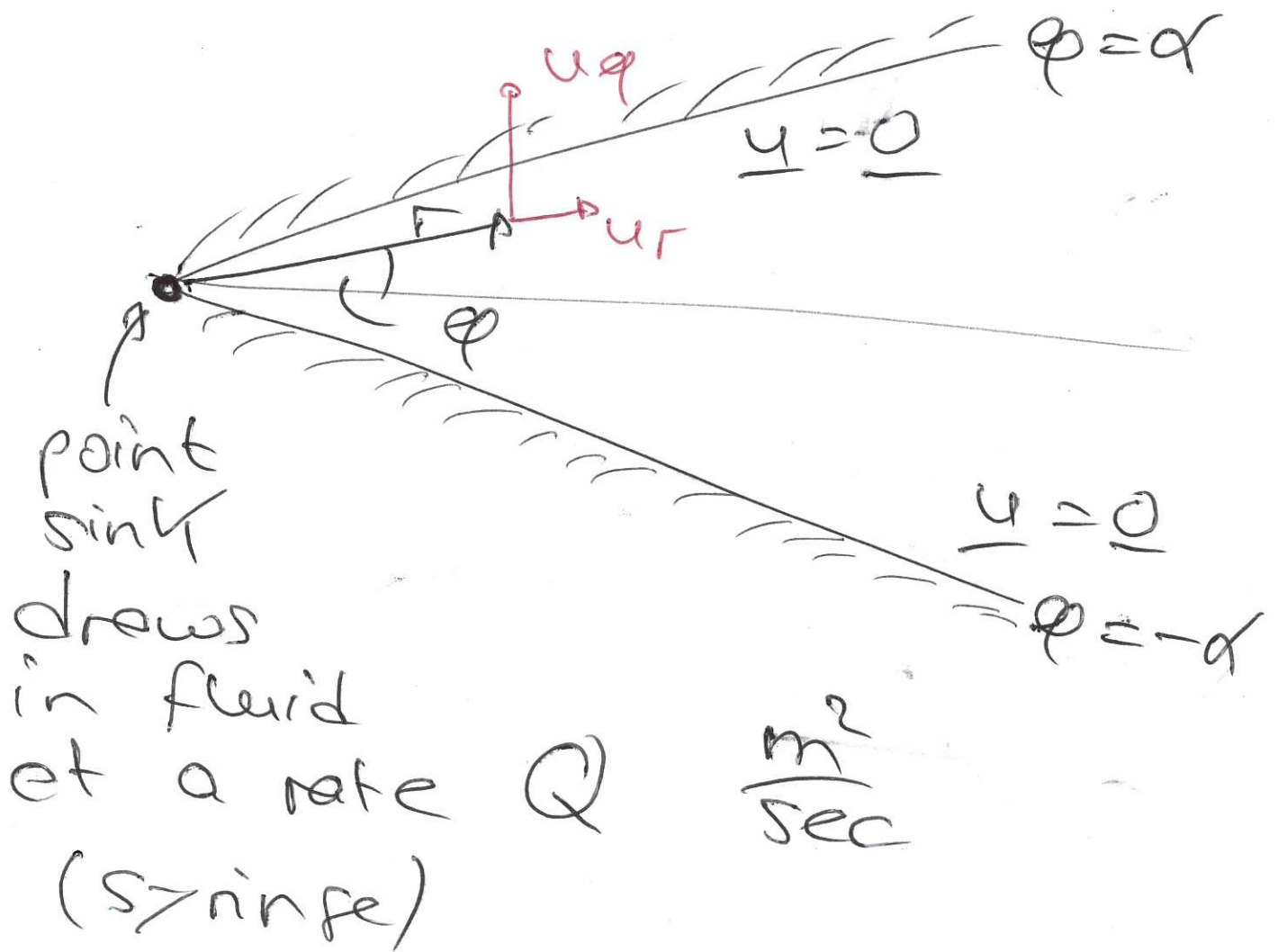
$\sqrt{\pi}$

So: $u = U F(\eta)$

$$u = \frac{U}{\sqrt{\pi \eta}} \int_0^{\infty} \exp\left(-\frac{1}{4} \xi^2\right) d\xi$$

where $\eta = \frac{y}{\sqrt{4 \nu t}}$

Example: Jeffrey-Hamel flow
 flow in wedge



$$\underline{u} = \underline{u}(\underline{r}, \varphi; \nu, Q, \alpha)$$

Dimensions:

(P)

f	$\frac{kg}{m^3}$
V	$\frac{m^2}{sec}$
Q	$\frac{m^2}{sec}$
α	1
φ	1
r	m
u	$\frac{m}{sec}$