

SG Perturbation methods (1)

- Problem has small parameters
- Problem is "easy" for $\varepsilon \rightarrow 0$

$$x^2 + \varepsilon x - 1 = 0$$

$$x(\varepsilon) = -\frac{1}{2}\varepsilon \pm \sqrt{\frac{1}{4}\varepsilon^2 + 1}$$

small ε : Taylor expand

$$x(\varepsilon) \approx \pm 1 - \frac{1}{2}\varepsilon \pm \frac{1}{8}\varepsilon^2 + \dots$$

Ansatz:

$$x(\varepsilon) = x_0 + \varepsilon x_1 + \varepsilon^2 x_2 + \dots$$

into eqn & collect powers of ε

$$(x_0^2 - 1) + \varepsilon(2x_0x_1 + x_0) + \varepsilon^2(x_1^2 + 2x_0x_2 + x_1) + \dots = 0$$

$$x(\varepsilon) \approx \pm 1 - \frac{1}{2}\varepsilon \pm \frac{1}{8}\varepsilon^2 + \dots$$

Note the structure of
the eqns:

(2)

- Lowest-order eqn is the full eqn for $\epsilon = 0$, which we assumed was "easy".
- Higher-order eqns provide corrections via a hierarchy of problems.
- Eqn. itself is satisfied ~~eqn~~ up to a quantifiable error.

Questions:

- How do we "guess" the right form of the expansion?
- Will it always work?

An ODE example

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Damped oscillator with weak damping.

$$\ddot{x} + \varepsilon \dot{x} + x = 0$$

$$x(t=0) = 1$$

$$\dot{x}(t=0) = 0$$

$$\Rightarrow x(t; \varepsilon)$$

Assume we don't know the form. Note: for $\varepsilon=0$: $x(t, \varepsilon=0) = \cos t$

Ansatz for small ε :

$$x(t; \varepsilon) = x_0(t) + \varepsilon x_1(t) + \varepsilon^2 x_2(t) + \dots$$

into problem:

$$\dot{x}(t) = \dot{x}_0(t) + \varepsilon \dot{x}_1(t) + \varepsilon^2 \dot{x}_2(t) + \dots$$

$$\ddot{x}(t) = \ddot{x}_0(t) + \varepsilon \ddot{x}_1(t) + \varepsilon^2 \ddot{x}_2(t) + \dots$$

into ODE

$$\underbrace{\ddot{x}_0 + \epsilon \ddot{x}_1 + \epsilon^2 \ddot{x}_2 + \dots}_\ddot{x} +$$

$$\underbrace{\dot{x}_0 + \epsilon \dot{x}_1 + \epsilon^2 \dot{x}_2 + \dots}_\dot{x} +$$

$$\underbrace{x_0 + \epsilon x_1 + \epsilon^2 x_2 + \dots}_x = 0$$

Collect powers of ϵ :

$$\left(\ddot{x}_0 + \dot{x}_0 \right) + \epsilon \left(\ddot{x}_1 + \dot{x}_0 + \dot{x}_1 \right) + \epsilon^2 \left(\ddot{x}_2 + \dot{x}_1 + \dot{x}_2 \right) + \dots = 0 \quad (1)$$

IC: $x(t=0) = 1$

$$\underbrace{x_0(t=0)} + \epsilon \underbrace{x_1(t=0)} + \epsilon^2 \underbrace{x_2(t=0)} + \dots = 1$$

$$\left(x_0(t=0) - 1 \right) + \epsilon \left(x_1(t=0) \right) + \epsilon^2 \left(x_2(t=0) \right) + \dots = 0 \quad (2)$$

$$\dot{X}(t=0) = 0$$

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$$\dot{X}_0(t=0) + \varepsilon \dot{X}_1(t=0) + \varepsilon^2 \dot{X}_2(t=0) + \dots = 0 \quad (3)$$

Now "solve" (1) - (3) by setting terms in increasing powers of ε to zero.

$O(\varepsilon^0)$:

$$\begin{aligned} \ddot{X}_0 + X_0 &= 0 \\ \dot{X}_0(t=0) &= 1 \\ X_0(t=0) &= 0 \end{aligned}$$

← inf-problem for $\varepsilon=0$

→ $X_0(t) = \cos(t)$

$O(\varepsilon^1)$:

$$\begin{aligned} \ddot{X}_1 + X_1 &= -\dot{X}_0 \\ X_1(t=0) &= 0 \\ \dot{X}_1(t=0) &= 0 \end{aligned}$$

→ $X_1(t)$

$O(\varepsilon^2)$:

$$\begin{aligned} \ddot{X}_2 + X_2 &= -\dot{X}_1 \\ X_2(t=0) &= 0 \\ \dot{X}_2(t=0) &= 0 \end{aligned}$$

⋮

Solve these problems
in sequence

(6)

$O(\varepsilon^0)$:

$$x_0(t) = \cos(t)$$

(This is the "easy" soln.
of the anf. problem
for $\varepsilon = 0$)

$O(\varepsilon^1)$:

$$\ddot{x}_1 + x_1 = -x_0 = \sin(t)$$

$$x_{1H} = A \cos t + B \sin t$$

$$x_{1P} = C t \sin t + D t \cos t$$

⋮

$$x_1 = A \cos t + B \sin t - \frac{1}{2} t \cos t$$

Apply ICS

$$x_1(t) = \frac{1}{2} \sin(t) - \frac{1}{2} t \cos t$$

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$O(\varepsilon^2)$: $\ddot{x}_2 + x_2 = -\dot{x}_1 = -\frac{1}{2}t \sin t$

$$x_{2H} = A \cos t + B \sin t$$

$$x_{2P} = -\frac{1}{8}t \sin t + \frac{1}{8}t^2 \cos t$$

x_{2P} satisfies I.C. too.

$$x_2(t) = x_{2P}(t)$$

⋮

$$x(t; \varepsilon) = \cos(t) + \varepsilon \left(\frac{1}{2} \sin(t) - \frac{1}{2}t \cos(t) \right) + \varepsilon^2 \left(-\frac{1}{8}t \sin(t) + \frac{1}{8}t^2 \cos(t) \right) + \dots$$

Mechanical oscillator with weak damping

- Governing (linear) ODE:

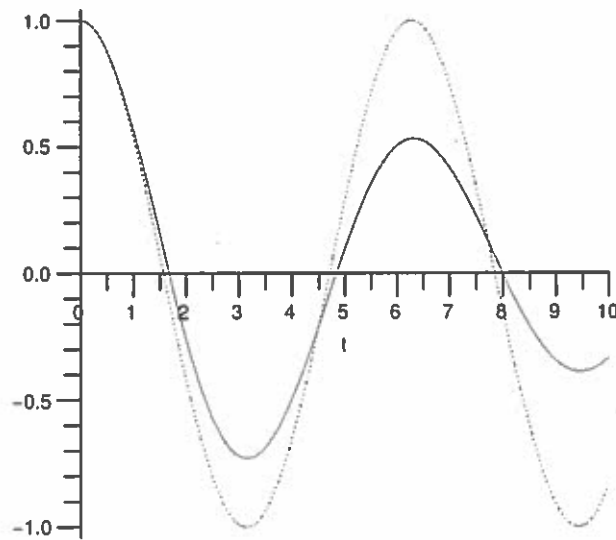
$$\ddot{x} + \epsilon \dot{x} + x = 0$$

subject to the initial conditions

$$x(t=0) = 1 \quad \text{and} \quad \dot{x}(t=0) = 0.$$

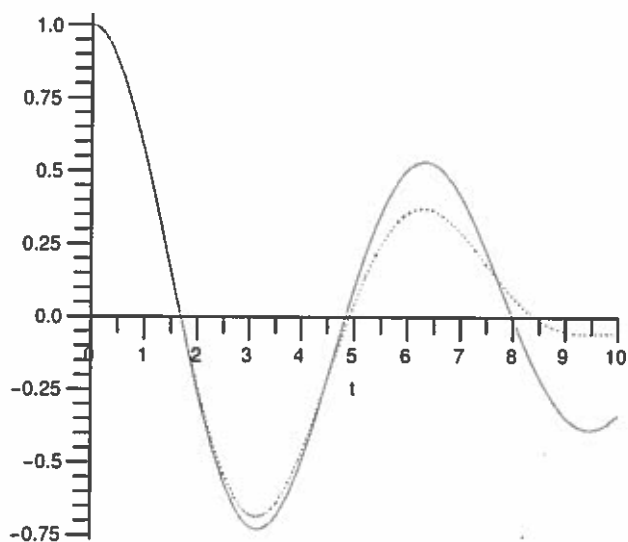
Comparison between perturbation solution and exact solution for $\epsilon = 0.2$

- One-term perturbation solution (red), exact solution (green):

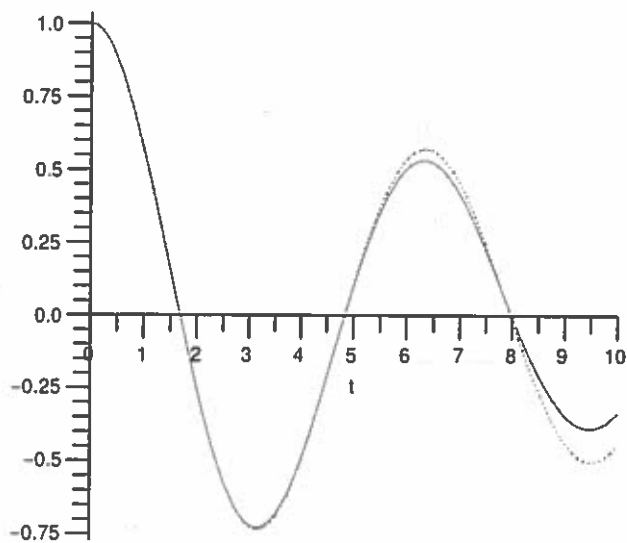


Comparison between perturbation solution and exact solution for $\epsilon = 0.2$

- Two-term perturbation solution (red), exact solution (green):

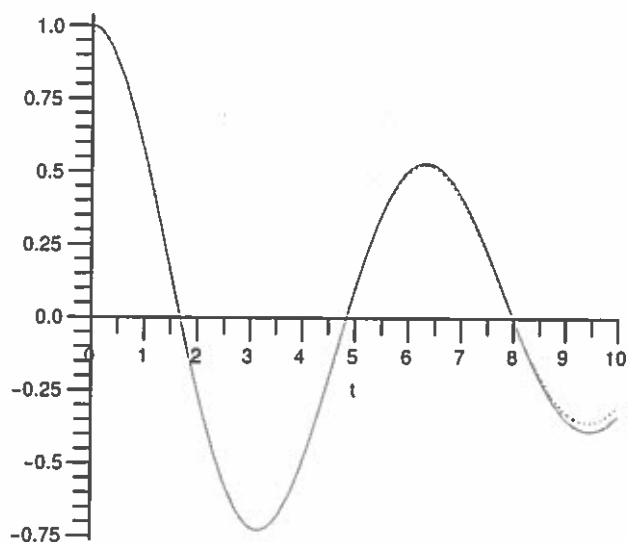


- Three-term perturbation solution (red), exact solution (green):

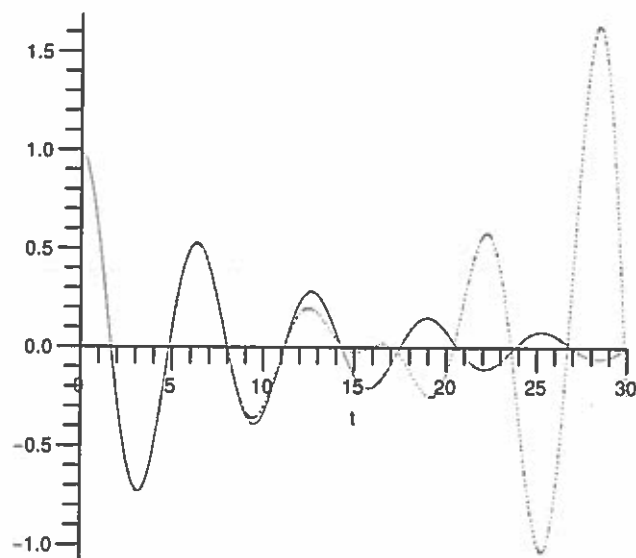


Comparison between perturbation solution and exact solution for $\epsilon = 0.2$

- Four-term perturbation solution (red), exact solution (green):



- Agreement over a finite time-interval is very pleasing. However, over sufficiently large times, the perturbation solution diverges:



observations:

① Perturbation soln. converges rapidly (at fixed time) as

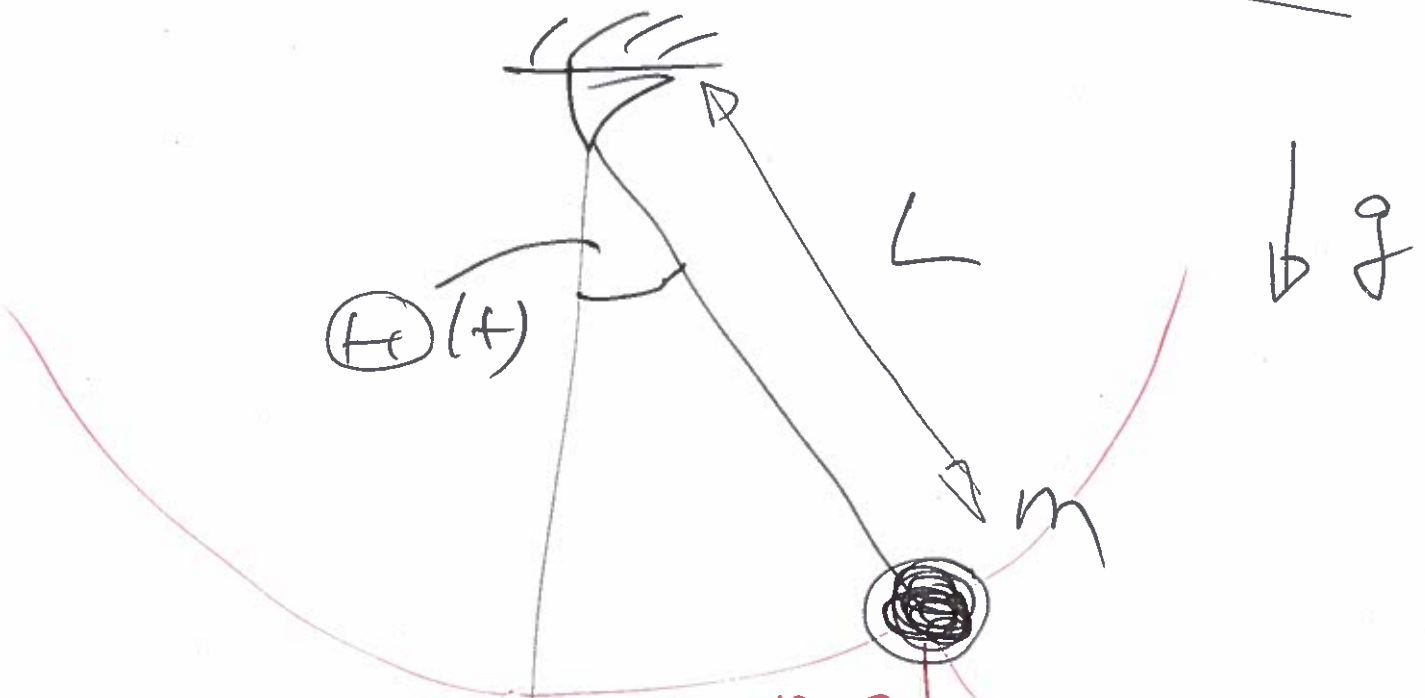
- the number of terms is increased
- ϵ is decreased (not shown)

② As $t \rightarrow \infty$ the solution becomes inaccurate

- because perturbation soln. contains t^n terms with $n > 0$
 - errors in the soln. of the ODE "accumulate".
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A nonlinear example

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Newton's law in tangential direction:



$$m \frac{d}{dt} \left(\underbrace{L \frac{d\theta}{dt}}_{\text{tangential velocity}} \right) = -mg \sin \theta$$

$$\ddot{\theta} + \omega^2 \sin \theta = 0$$

$$\omega^2 = \frac{g}{L}$$

IC:

$$h(t=0) = \varepsilon$$

$$\frac{dh}{dt} \Big|_{t=0} = 0$$

IC