MAGIC: Ergodic Theory Lecture 7 - Entropy

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commutes (up to sets of measure zero) and $\mu \circ \phi^{-1} = m$ (i.e. $\mu(\phi^{-1}B) = m(B) \ \forall B \in \mathcal{A}$). It is natural to look for invariants. To each mpt T we will associate a number - its *entropy* h(T). If S, T are isomorphic then h(S) = h(T). (Conversely, if $h(S) \neq h(T)$ then S, T cannot be isomorphic.) Throughout, $\log = \log_2$.

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Information and entropy of a partition Let (X, \mathcal{B}, μ) be a probability space.

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This motivates defining the 'information function' as

$$I(\alpha)(x) = \sum_{A \in \alpha} \chi_A(x)\phi(\mu(A))$$

for an appropriate choice of function ϕ .

Suppose α and β are two partitions.

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$$\alpha \lor \beta = \{ A \cap B \mid A \in \alpha, B \in \beta \}.$$

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 α,β are independent if $\mu(A \cap B) = \mu(A)\mu(B) \ \forall A \in \alpha, B \in \beta$.

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Definition

The information function of $\boldsymbol{\alpha}$ is

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The entropy of α is the average amount of information:

$$H(\alpha) = \int I(\alpha) d\mu = -\sum_{A \in \alpha} \mu(A) \log \mu(A).$$

Conditional information & entropy

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For example: if β is a partition then the set of all unions of elements of β is a σ -algebra (also denoted by β).

How much information do we gain by knowing which element of α we are in, given we know which element of β we are in?

Recall conditional expectation:

$$\mathbb{E}(\cdot \mid \mathcal{A}) : L^1(X, \mathcal{B}, \mu) \longrightarrow L^1(X, \mathcal{A}, \mu).$$

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 $\mathbb{E}(f \mid \mathcal{A})$ is determined by

1. $\mathbb{E}(f \mid A)$ is A-measurable,

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$$\int_{A} \mathbb{E}(f \mid A) d\mu = \int_{A} f d\mu \ \forall A \in A.$$

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 $\mathbb{E}(f|\mathcal{A})$ is the best \mathcal{A} -measurable approximation to f.

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$$\mathbb{E}(f \mid \beta)(x) = \sum_{B \in \beta} \chi_B(x) \frac{\int_B f \, d\mu}{\mu(B)}.$$

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Let $A \in \mathcal{B}$. The conditional probability of A given a sub- σ -algebra \mathcal{A} is

$$\mu(\mathcal{A}|\mathcal{A}) = \mathbb{E}(\chi_{\mathcal{A}}|\mathcal{A}).$$

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Consider the σ -algebra β generated by a partition β .



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Note:

$$\mu(A \mid \beta) = \sum_{B \in \beta} \chi_B(x) \frac{\mu(A \cap B)}{\mu(B)}.$$

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The basic identities:

$$I(\alpha \lor \beta \mid \gamma) = I(\alpha \mid \gamma) + I(\beta \mid \alpha \lor \gamma)$$

$$H(\alpha \lor \beta \mid \gamma) = H(\alpha \mid \gamma) + H(\beta \mid \alpha \lor \gamma)$$

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$$\beta \leq \gamma \implies I(\alpha \lor \beta \mid \gamma) = I(\alpha \mid \gamma)$$

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follow from the basic identities, 3 follows from Jensen's ineq.

Entropy of an mpt relative to a partition

We can now start to define the entropy h(T) of an mpt T. We first define the entropy of T relative to a partition. We need the following:

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Subadditive lemma

Suppose $a_n \in \mathbb{R}$ is subadditive: $a_{n+m} \leq a_n + a_m$.

Then
$$\lim_{n\to\infty} \frac{a_n}{n}$$
 exists and equals $\inf_n \frac{a_n}{n}$ (could be $-\infty$).

Let $T : (X, \mathcal{B}, \mu) \to (X, \mathcal{B}, \mu)$ be an mpt. Let α be a finite or countable partition.

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$$H(T^{-1}\alpha) = -\sum_{A \in \alpha} \mu(T^{-1}A) \log \mu(T^{-1}A)$$
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Define

$$H_n(\alpha) = H\left(\bigvee_{j=0}^{n-1} T^{-j}\alpha\right).$$

$$H_{n+m}(\alpha) = H\left(\bigvee_{j=0}^{n+m-1} T^{-j}\alpha\right)$$

$$\begin{aligned} H_{n+m}(\alpha) &= H\left(\bigvee_{j=0}^{n+m-1}T^{-j}\alpha\right) \\ &\stackrel{\text{basic}}{=} H\left(\bigvee_{j=0}^{n-1}T^{-j}\alpha\right) + H\left(\bigvee_{j=n}^{n+m-1}T^{-j}\alpha\middle|\bigvee_{j=0}^{n-1}T^{-j}\alpha\right) \end{aligned}$$

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Hence $H_n(\alpha)$ is subadditive.

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'Entropy = average amount of information from the present, given the past'

Entropy of an mpt

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Potential problem: working from the definitions, this quantity seems impossible to calculate!

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Definition

A finite or countable partition α is a *generator* for ${\cal T}$ if ${\cal T}$ is invertible and

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(i.e. \mathcal{B} is the smallest σ -algebra that contains all elements of all the partitions $\bigvee_{j=-(n-1)}^{n-1} T^{-j} \alpha$). We say that α is a *strong generator* if

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Remark

To check whether a partition α is a strong generator (resp. generator) it is sufficient to show it separates μ -a.e. pair of points:

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Recall:

$$h_{\mu}(T) = \sup h_{\mu}(T, \alpha)$$

where the supremum is taken over all partitions of finite entropy.

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Sinai's theorem tells us that this supremum is acheived when α is a generator or a strong generator.

Let α be a finite or countable partition with $H(\alpha) < \infty$.

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This allows us to calculate the entropy of many of our favourite examples.

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Recall that the Markov measure μ_P is defined on cylinder sets by:

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Let $\alpha = \{[1], \ldots, [k]\}$ denote the partition of Σ_A^+ into cylinders of length 1.



Easy check: $H(\alpha) < \infty$ Easy check: $\alpha_n = \bigvee_{j=0}^{n-1} \sigma^{-j} \alpha = \{[i_o, \dots, i_{n-1}]\}\$ = the partition of Σ_A^+ into cylinders of length n

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= $-\sum_{i_o,...,i_{n-1}} p_{i_0} P_{i_0 i_1} \dots P_{i_{n-2} i_{n-1}} \log(p_{i_0} P_{i_0 i_1} \dots P_{i_{n-2} i_{n-1}})$

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Easy check:
$$H(\alpha) < \infty$$

Easy check: $\alpha_n = \bigvee_{j=0}^{n-1} \sigma^{-j} \alpha = \{[i_o, \dots, i_{n-1}]\}$
= the partition of Σ_A^+
into cylinders of length n

Hence α is a strong generator, as α_n separates points. Hence we can apply Sinai's theorem:

$$\mathcal{H}\left(\bigvee_{j=o}^{n-1} \sigma^{-j} \alpha\right)$$

$$= -\sum_{i_{o},...,i_{n-1}} \mu[i_{0},...,i_{n-1}] \log \mu[i_{0},...,i_{n-1}]$$

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$$= -\sum_{i_{o},...,i_{n-1}} p_{i_{0}} \log p_{i_{0}} \dots (n-1) \sum_{i,j} p_{i_{j}} \log P_{i_{j}}.$$

$$h_{\mu}(\sigma) \stackrel{\text{Sinai}}{=} h_{\mu}(\sigma, \alpha)$$

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$$h_{\mu}(\sigma) \stackrel{\text{Sinai}}{=} h_{\mu}(\sigma, \alpha)$$
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$$-\sum_{i,j} p_i P_{ij} \log P_{ij}.$$

Remark

If μ is the Bernoulli- (p_1, \ldots, p_k) measure then

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This suggests that there is a lot of redundancy in English (good for error-correcting!). See Shannon's book on Information Theory.

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Theorem

If T, S are isomorphic then $h_{\mu}(T) = h_m(S)$.

Let α be a finite or countable partition of Y with $H_m(\alpha) < \infty$.

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Hence $h_{\mu}(T, \phi^{-1}\alpha) = h_m(S, \alpha)$.

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- ▶ φ is a bijection, except on the countable set of points which have non-unique base 2 expansions,
- λ = μφ⁻¹ (clear on dyadic intervals, follows for all sets by the Kolmogorov Extension Theorem).

Hence $Tx = 2x \mod 1$ with Lebesgue measure λ and the full one-sided 2-shift σ with the Bernoulli-(1/2, 1/2) measure μ are isomorphic.

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$$h_{\lambda}(T) = \log 2 = h_{\mu}(\sigma).$$

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2-sided aperiodic Markov shifts with the same entropy are isomorphic.

(The one-sided case is far more subtle.)

Being isomorphic to a Bernoulli shift is a useful and desirable property for a mpt to possess.

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In general, a mpt that exhibits some form of 'hyperbolicity' is, when equipped with a suitable measure, Bernoulli.

For example, hyperbolic toral automorphisms are Bernoulli.

Next lecture

Entropy has been defined in a purely measure-theoretic setting.

There is a topological analogue in the setting of continuous transformations of compact metric spaces: topological entropy.

We will define this and study the connections between measure-theoretic and topological entropy.